



Modeling black hole mergers in long gamma ray bursts

Agnieszka Janiuk

Center for Theoretical Physics,
Polish Academy of Sciences, Warsaw



www.cft.edu.pl

Collaboration:

Szymon Charzyński (Cardinal
Wyszyński University, Dept. of
Mathematics),



Michał Bejger (Copernicus
Astronomical Center)



**Recent paper: [ArXiv:1310.4869](https://arxiv.org/abs/1310.4869)
(Astronomy and Astrophysics in press)**

Gamma Ray Bursts

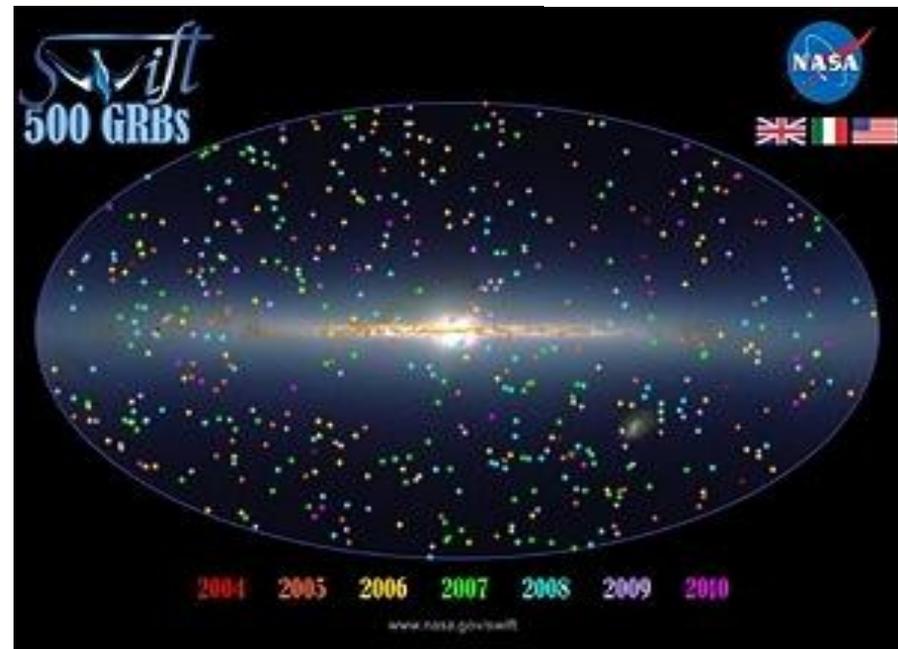
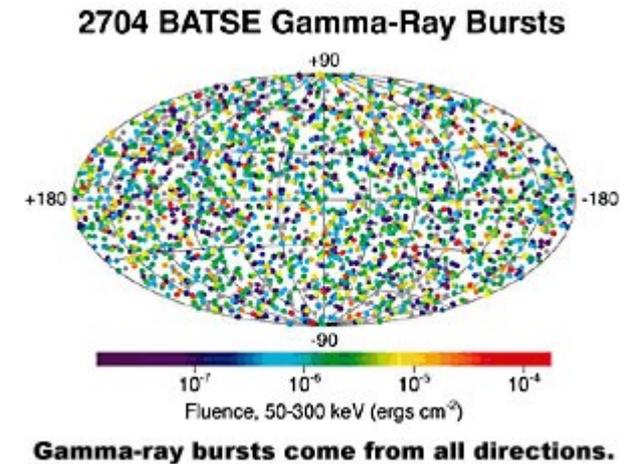
First detected in 1967

Isotropically distributed

At cosmological distances

Group in Long-Soft and Short-Hard

Prompt emission in Gamma plus afterglow up to IR/Radio bands

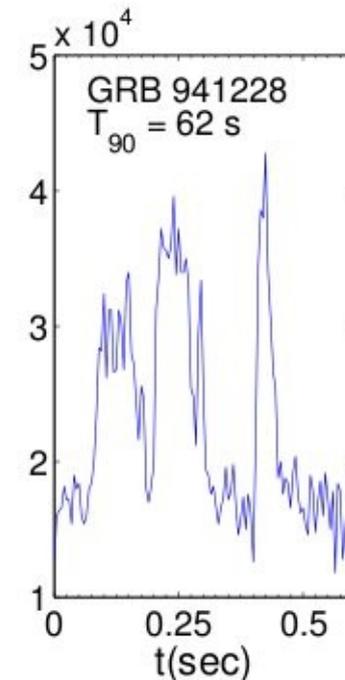
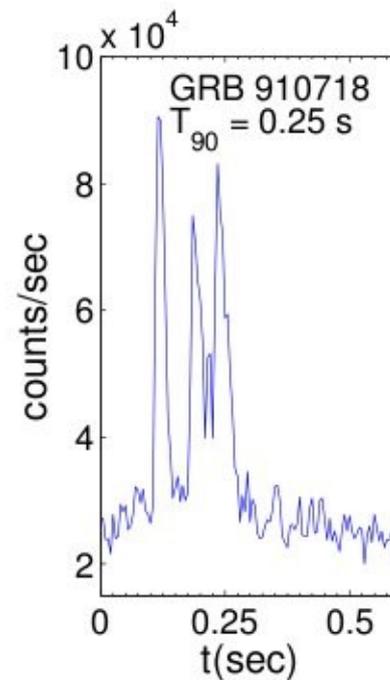


This all-sky map shows the locations of Swift's 500 gamma-ray bursts, color coded by the year in which they occurred. In the background, an infrared image shows the location of our galaxy and its largest satellites. Credit: NASA/Swift/Francis Reddy.

Time profiles of GRBs:

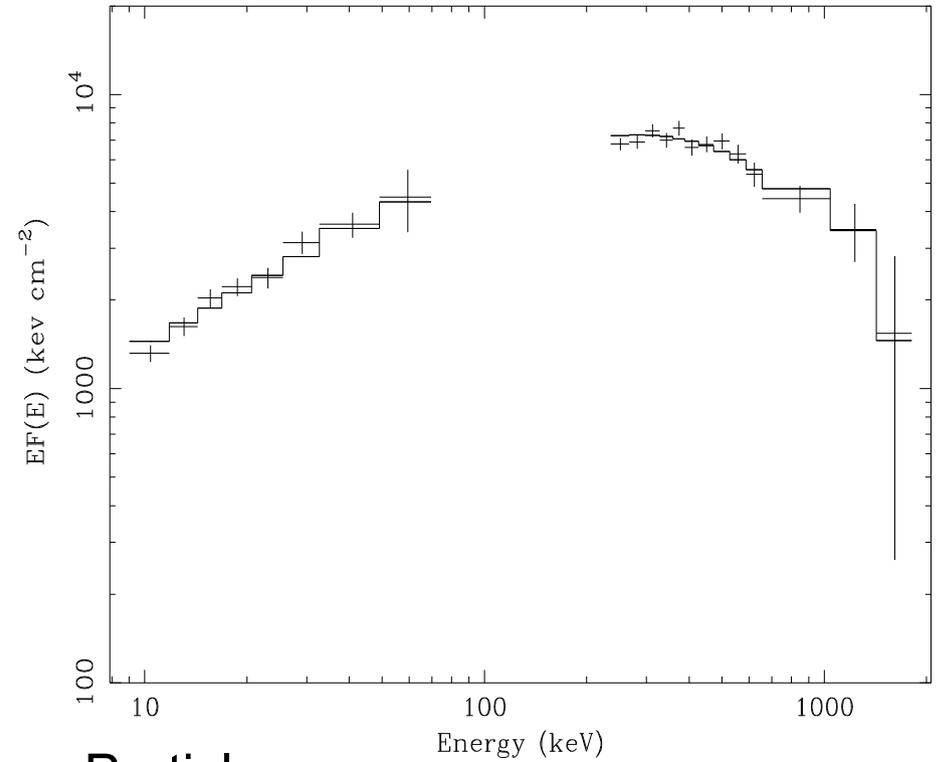
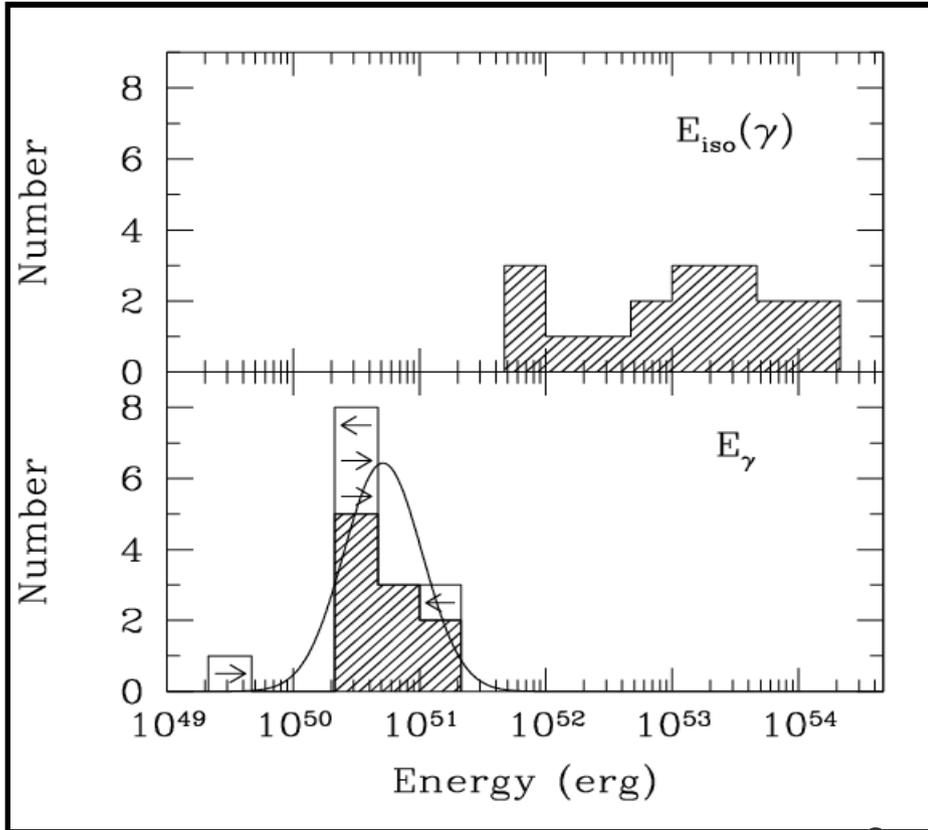
Duration given
by the T_{90}
estimate

Fast variability



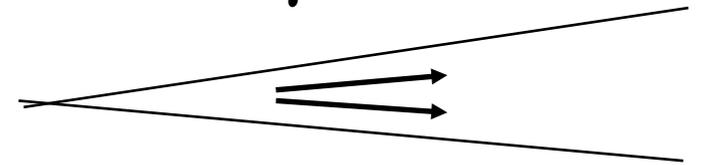
Energetics of GRBs

GRB 990510



Particles
move in a
cone

$$\gamma = \theta^{-1}$$



$$E_{tot} = \epsilon_\gamma^{-1} E_\gamma = \epsilon_\gamma^{-1} \frac{\theta^2}{2} E_{\gamma iso}$$

Source of energy for GRBs

Required energy: $\sim 10^{52}$ erg
(to be radiated)

$$E = G M^2/r$$

Typical models of GRBs
postulate the gravitational
potential energy

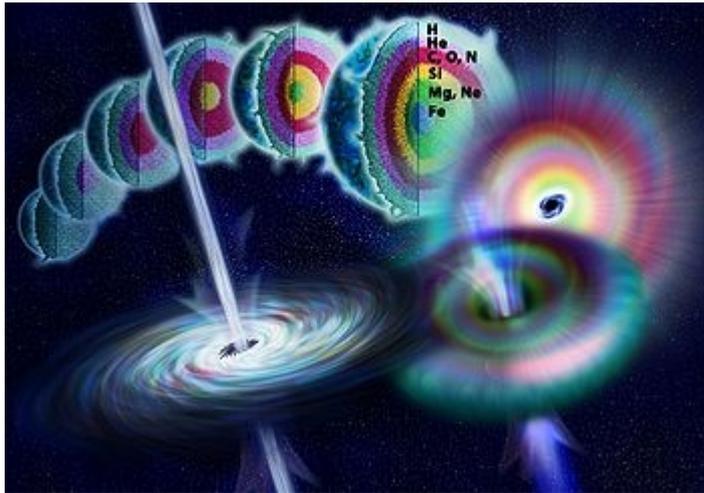
$$M = 1-10 \text{ Solar Masses}$$
$$r = 3-10 \text{ km}$$

The most efficient mechanism
for energy production is accretion
of matter onto a compact star

$$E = 10^{53}-10^{54} \text{ erg}$$

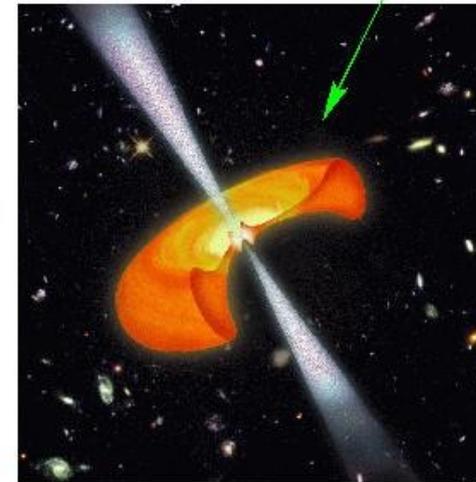
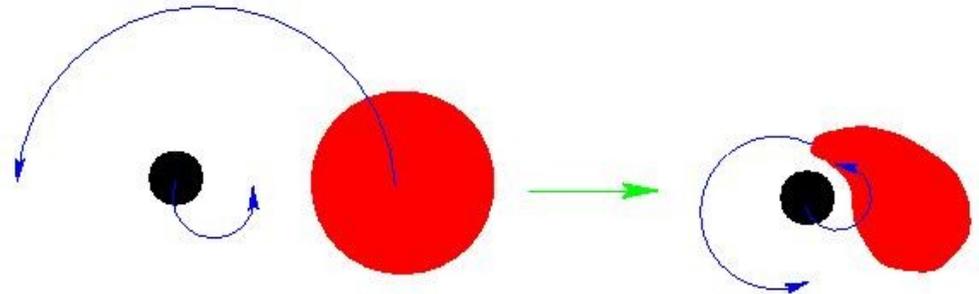
at efficiency of 1-10%

Two standard types of models



Massive star explosion and mass fallback from the envelope

Paczyński (1998); McFadyen & Woosley (1999)



Compact binary merger: neutron star disruption

Eichler et al. (1989); Ruffert & Janka (1999)

Another possible progenitor: high mass X-ray binary

Close binary: massive OB star plus compact remnant

Massive star is spun up by the interaction in binary system

Examples: Cyg X-3, IC 10 X-1, NGC 300 X-1

Ultimately, the compact star enters WR's envelope

Tidal squeezing triggers the core collapse in the primary

Two black holes merge

Matter accretes on the product

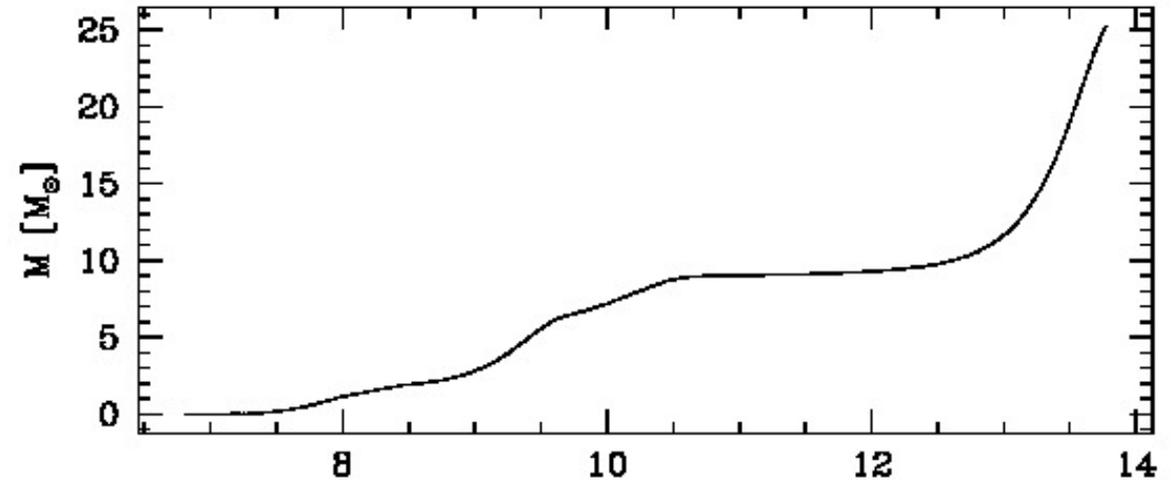


Zhang & Fryer (2001); Barkov & Komissarov (2010); Belczynski et al. (2008; 2013); Church et al. (2012)

Model of pre-supernova star

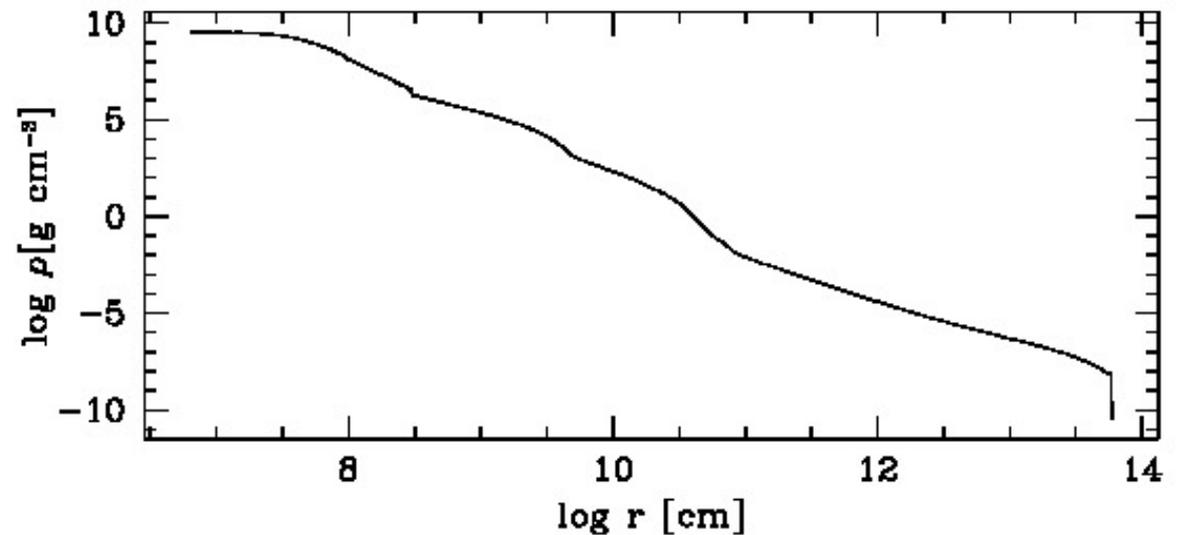
Pre-supernova star
(Woosley & Weaver 1995)

Enclosed mass of $25 M_{\text{Sun}}$



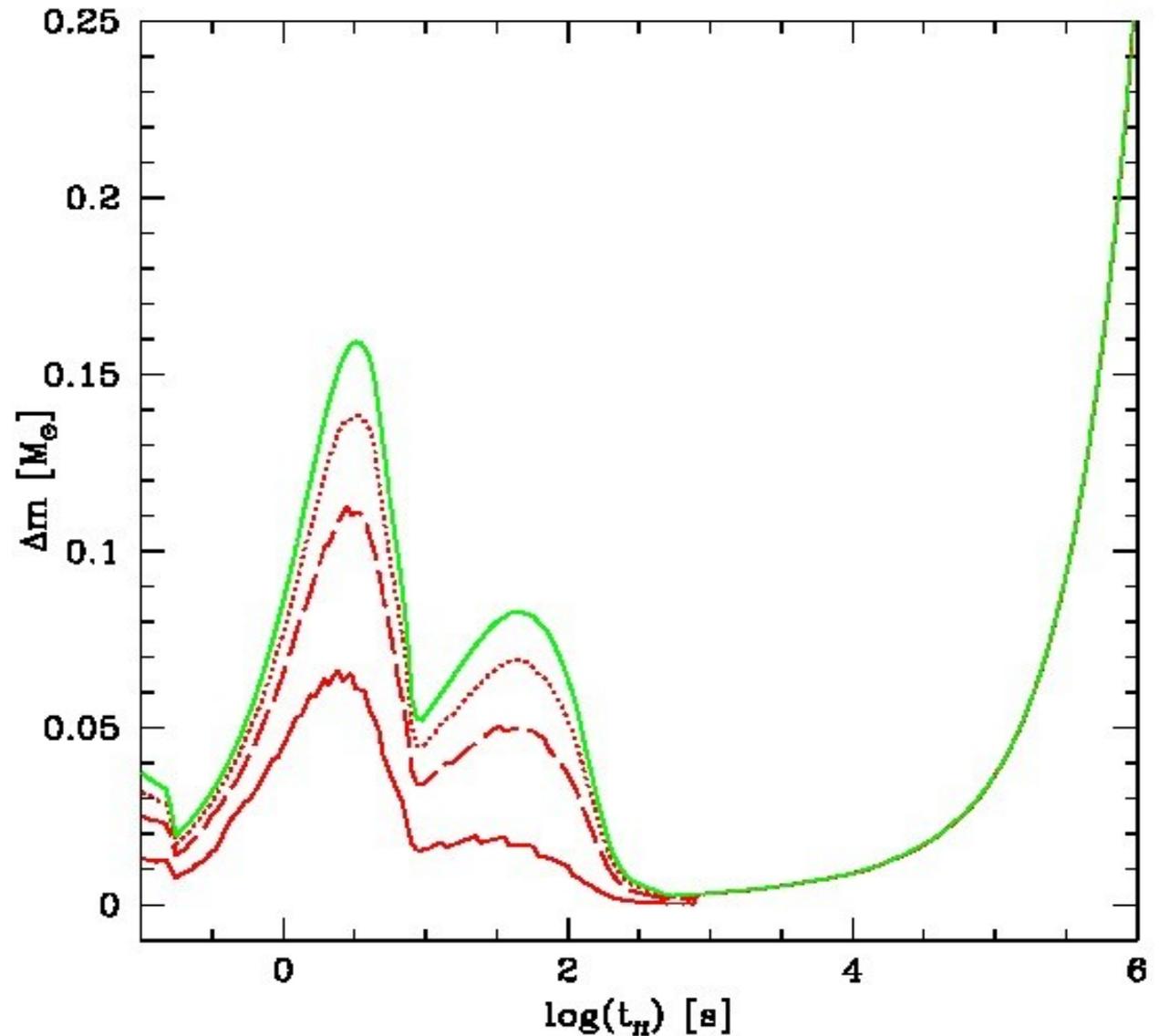
Density distribution:
chemical composition of an
evolved star (Fe, Si, C, O,
He, H)

Iron core of mass $1.4 M_{\text{Sun}}$



Mass of the envelope shell during the collapse (green:total; red: contained in the rotating torus).

The three lines show the models with various normalisation of the specific angular momentum in the envelope: $x=1.5$ (solid), $x=3.0$ (dashed) and $x=7.0$ (dotted).



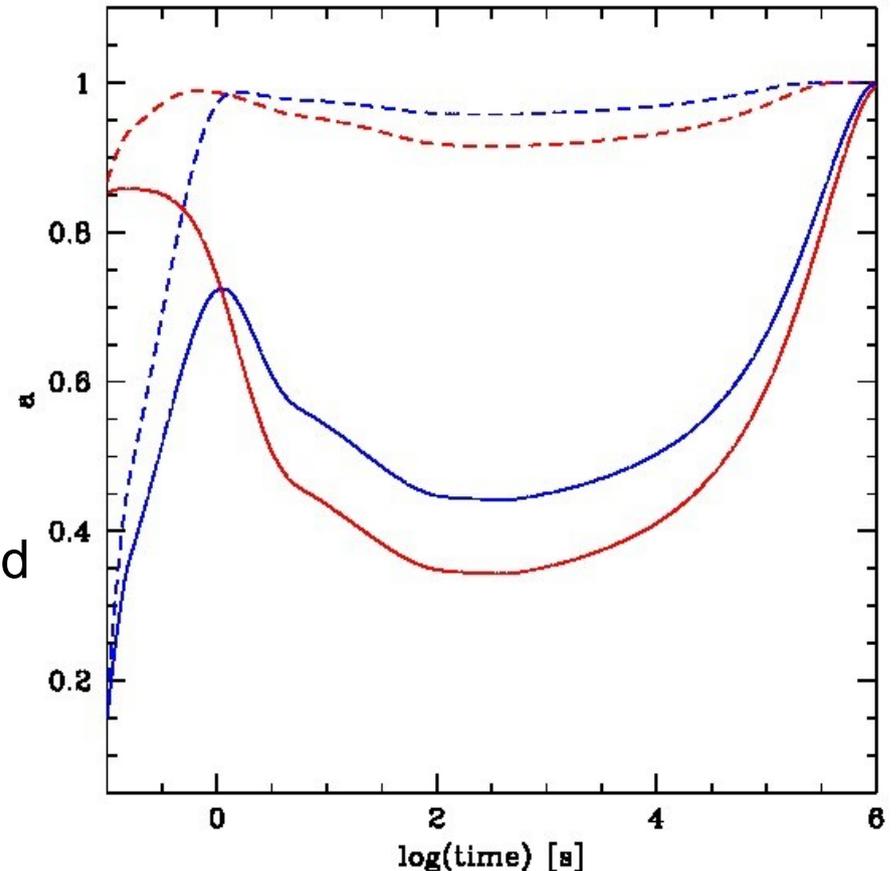
Spinning up the envelope and core black hole

We adopt specific angular momentum distribution in the star (differential rotation)

$$l_{spec}(r, \theta) = l_0 (1 - |\cos(\theta)|)$$

The infalling envelope matter adds mass and spins the black hole. The rotationally supported torus must obey (Bardeen et al. 1972):

$$l_{spec} > l_{crit} = \frac{2G M_{BH}}{c} \sqrt{2 - a + 2\sqrt{1 - a}}$$



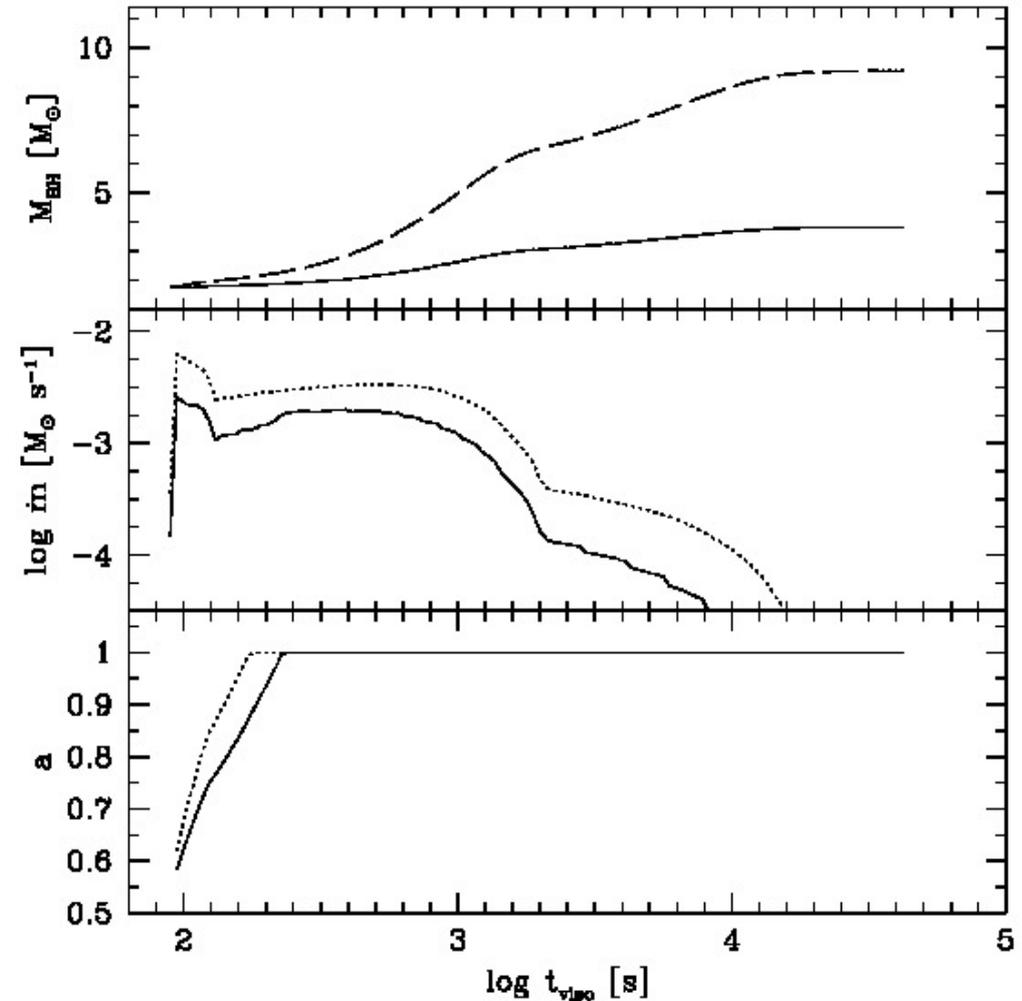
Janiuk & Proga (2008); Janiuk,
Moderski & Proga (2008)

Spinning up the envelope and core black hole

The companion BH transfers its orbital specific angular momentum

$$\Delta l = \frac{dJ_2}{dM} = \frac{M_2}{2} \sqrt{\frac{Gr}{M(r)}} \left(1 + \ln \frac{r_2}{r}\right)$$

Torus accretion: BH spins up to maximum rate. Some (most ?) of its mass will be lost in wind.



McKinney et al. (2006); Kumar et al. (2008);
Janiuk et al. (2004; 2007; 2010; 2013)

Evolution scenarios

- We proposed two representative scenarios for the pre-merger configuration
 - Homologous accretion on the envelope, when both the rotating torus and material from the poles contribute to the growth of the primary black hole
 - Only torus accretion, while the material from the poles is expelled. Some torus mass might further be expelled through winds
- The primary BH grows in mass to about $\sim 3 - 9 M_{\text{Sun}}$. This phase may last up to ~ 500 s (jet emission). Then the secondary BH sinks to merge with it. The spin of the primary at merger time is either $\sim 0.7 - 1$.

Merging two black holes

- Numerics done with Cactus Computational Toolkit (Goodale et al. 2003; Loeffler et al. 2012)
- 3+1 split of Einstein equations; Cauchy initial value problem solved with BSSN method (Shibata & Nakamura 1995; Baumgarte & Shapiro 1999)
- 3D Cartesian grid; adaptive mesh, reflection symmetry assumed to reduce number of grid points and computer requirements



www.einsteintoolkit.org

*Interdisciplinary Center for
Mathematical and
Computational Modeling,
Warsaw University*

We track numerically
the very last stage of
BBH merger:

initial separation of
6M; inspiral, merger
and ringdown

Quasicircular orbits,
mass ratio $q=1-3$

Primary spin $a=0-0.9$,
directed
perpendicularly to
orbital plane

Secondary is spinless

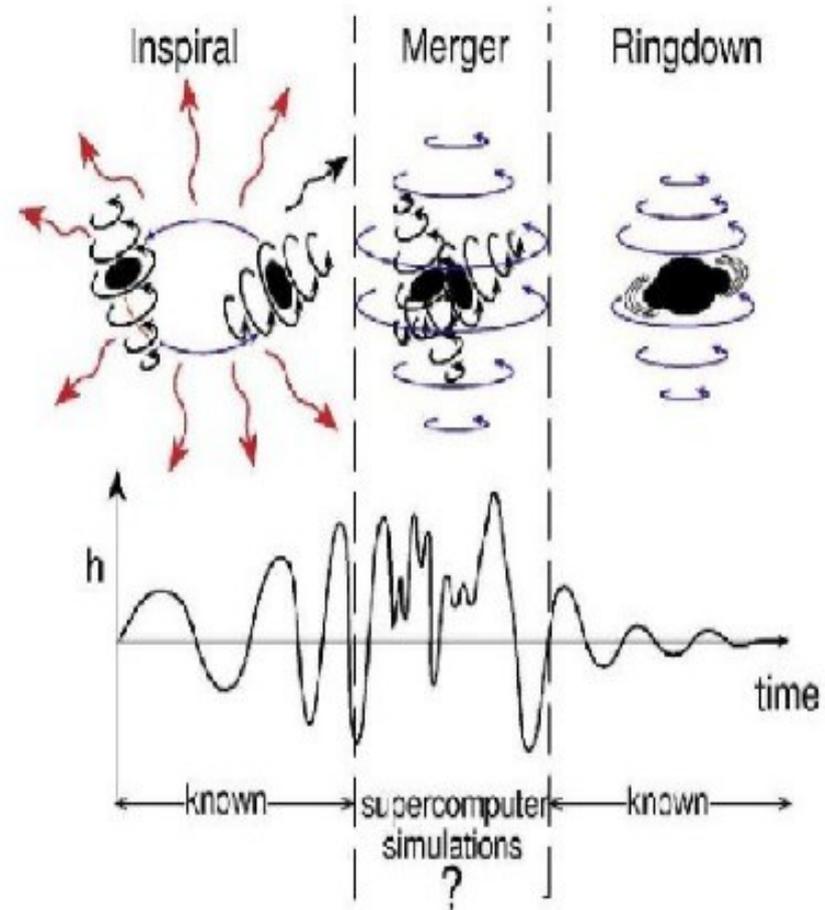


Fig. 1. The three phases of black hole merger (courtesy Kip Thorne).

Trajectories of the components during merger

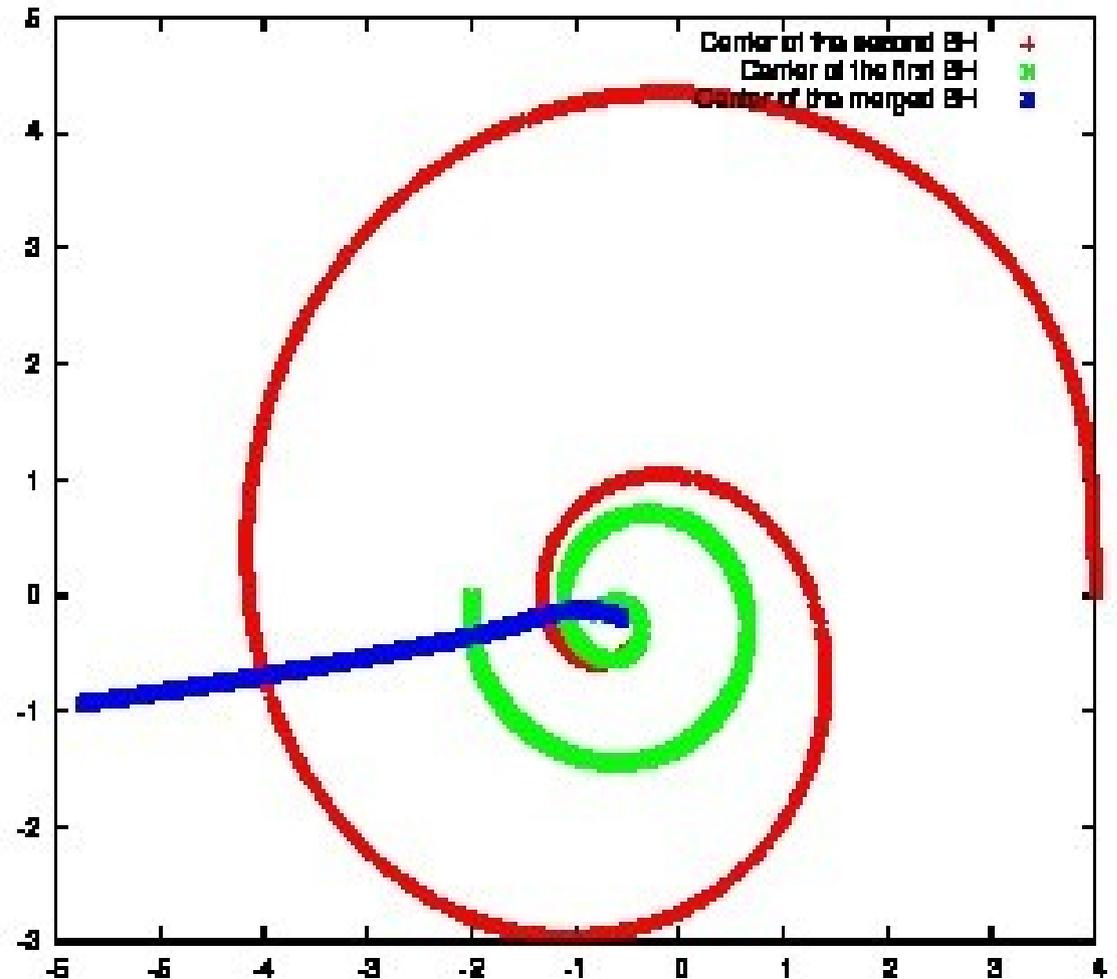
- Parameters:

$$m_1 = 0.632, m_2 = 0.316, s_1 = 0.9$$

- ADM mass ratio $M_1/M_2 = 3.0$

- Resulting final ADM mass and dimensionless spin

$$M_3 = 1.34, a_3 = 0.76$$



Gravitational recoil

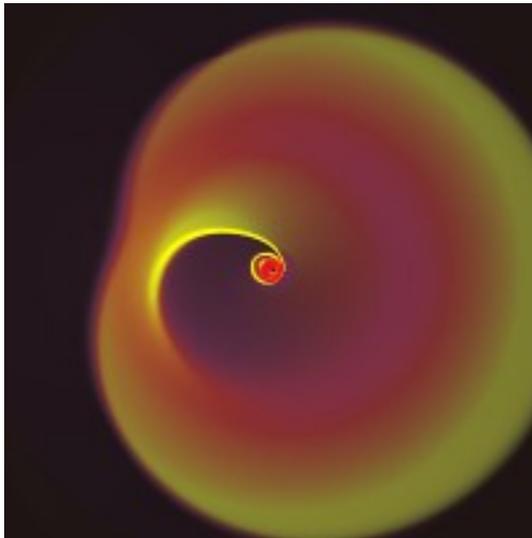
- Total linear momentum radiated from the system through gravitational waves is computed through the coefficients A^{lm} of multipole expansion of the Weyl scalar (Alcubierre 2008).
- Recoil vector remains in the orbital plane, because we assumed reflection symmetry; in general, it does not have to be the case
- We obtained velocity of the product, depending on spins and mass ratio of components, 200-300 km/s



Observational perspectives

Larger recoil speed for $q \sim 1$: if primary BH mass small due to wind taking most of envelope's, the offset of GRB afterglow is possible.

The merger product would leave the host galaxy if $v \sim 2000$ km/s. Possible if both BHs have extremely large spins (Tichy & Maronetti 2007).



Recoil kick directed into circumbinary disk plane can alter the distribution of specific angular momentum (Rossi et al. 2010).

If magnetic fields are involved, expansion of dual jets driven by generalized Blandford-Znajek mechanism (Palenzuela et al. 2010)

Summary

■ The long duration GRB may originate from a merger of a close compact binary system, containing a high mass evolved star and a black hole.

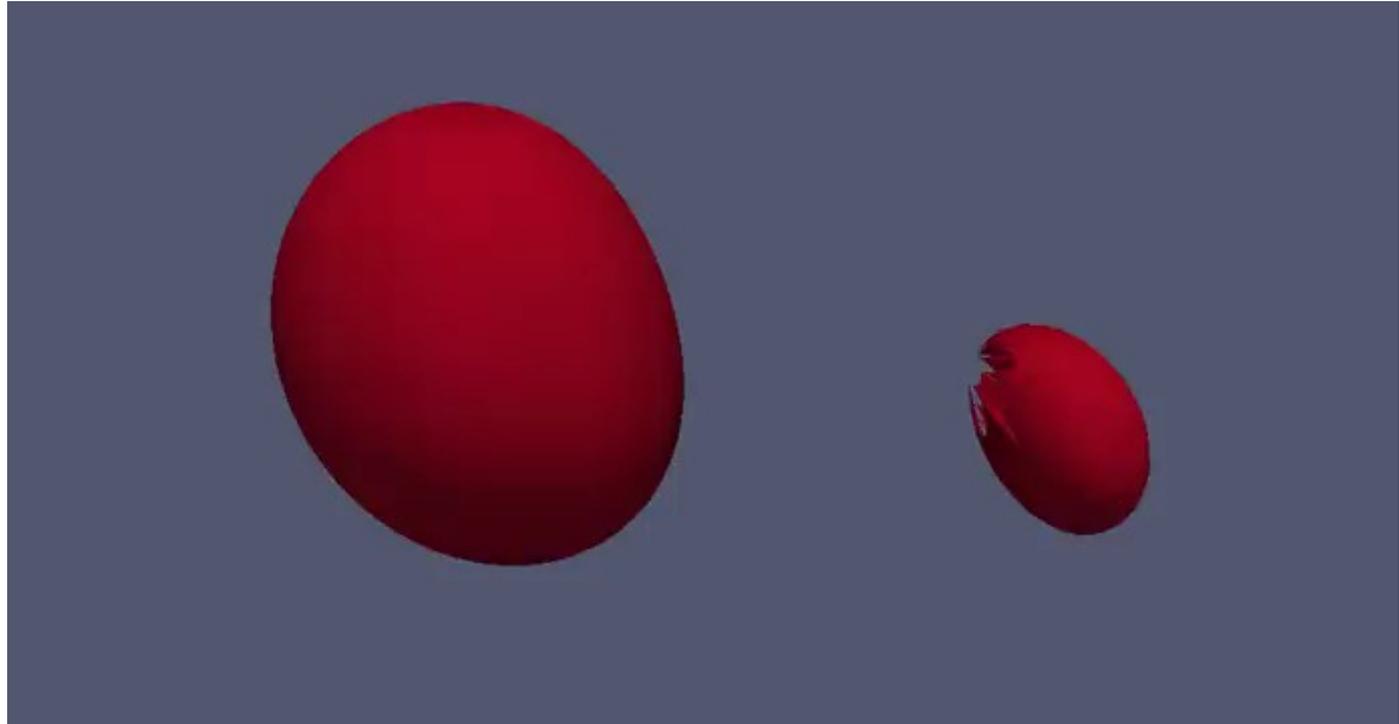
■ The event can be divided into stages:

1. The onset of the core collapse in the primary star, connected with the tidal interaction with secondary black hole. The inner shells of the envelope are spun up by the companion, and accrete onto the primary BH, increasing its mass and rotation spin
2. The merger of two black holes, surrounded by a circumbinary torus; gravitational waves; kick
3. Accretion of a remnant mass onto the BH merger product

■ Possible observational consequences

- Electromagnetic signal from the jets, POSSIBLY DOUBLE
- Gravitational waves
- Possible delay and offset of the secondary GRB signal / afterglow emission, due to recoil

Merging event horizons



Szilagyi et al. (2009, Phys Rev). Animation illustrates computations with the pseudo-spectral method, mass ratio 2, spins 0.4. Source: www.black-holes.org